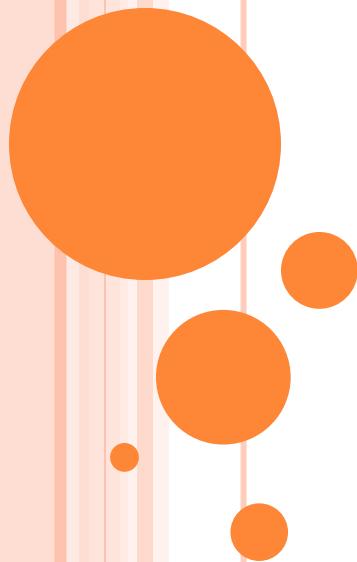


DISCOVERING HIGGS IN SUSY GUTS WITH TAU LEPTONS AT LHC

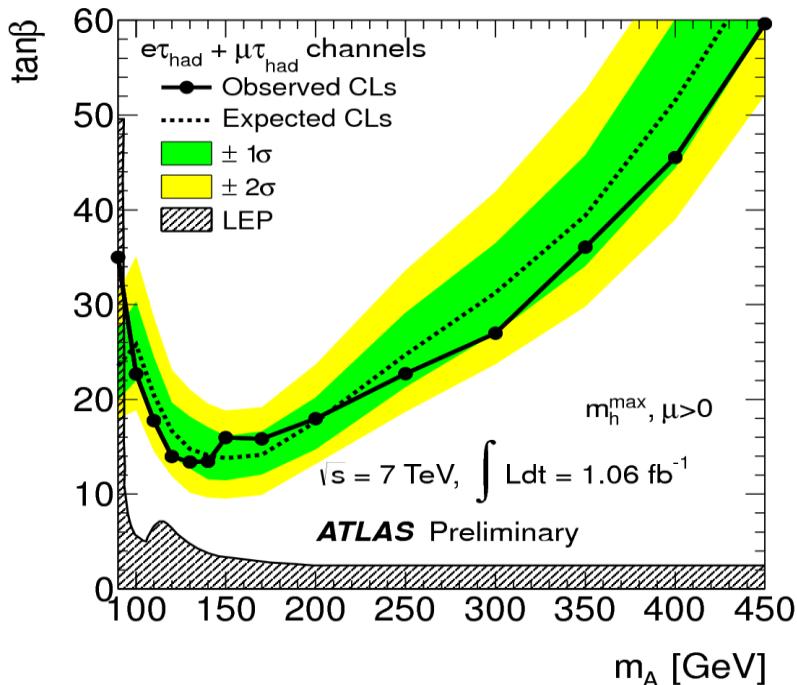
Kesheng Yang

(In collaboration with Prof. C. Kao)

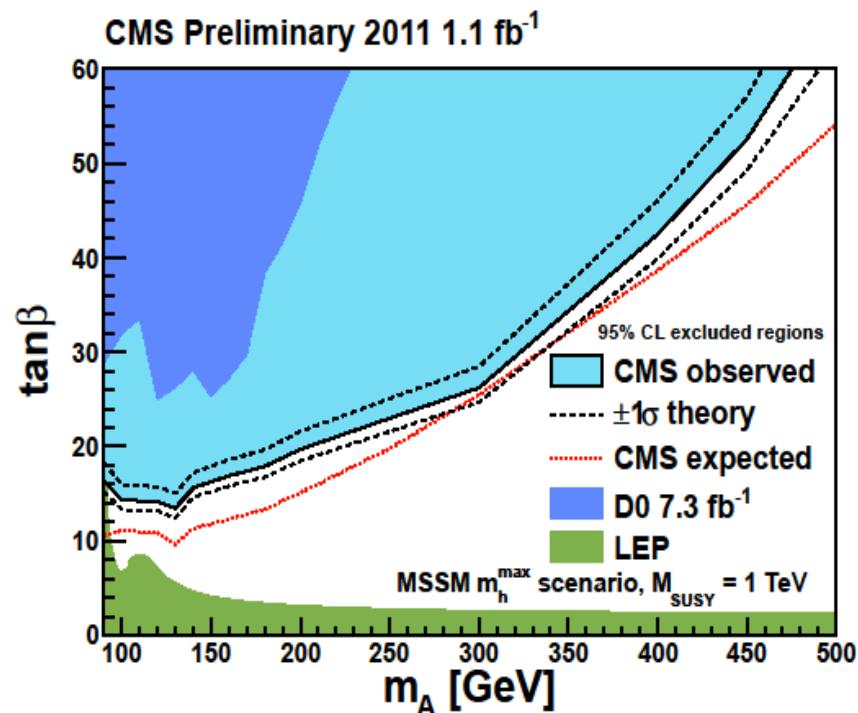
Homer L. Dodge Department of Physics
University of Oklahoma



MSSM Higgs searches from ATLAS and CMS



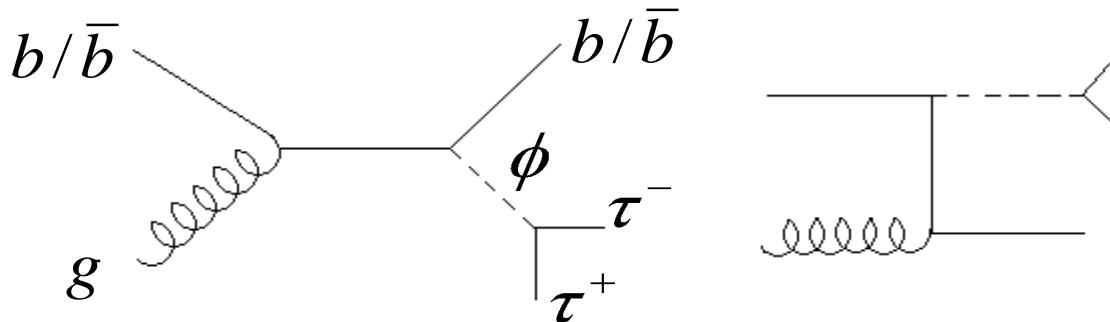
(ATLAS-CONF-2011-132)



(CMS-PAS-HIG-11-009)

Higgs Searches in the MSSM

- Higgs production associated with one b jet, and Higgs decay into $\tau^+\tau^-$ pairs.



- One b jet in the final state is helpful to handle the fake jets by applying b tagging technique.
- Neutral scalar ϕ could be h , H or A .
 - ❖ In low mass region, $m_h \sim m_A$. If $|m_h - m_A| < 10\% m_A$, we add the contribution from h together with that from pseudo scalar.
 - ❖ In high mass region, $m_H \sim m_A$, If $|m_H - m_A| < 10\% m_A$, we put the contribution from h together with that from pseudo scalar.
- Scalars predominantly decay to $b\bar{b}$ ($\sim 90\%$), and $\tau^+\tau^-$ ($\sim 10\%$).

- The biggest decay mode of tau pairs is one into hadronic jet with the other into electron or muon.

$$BF(\tau \rightarrow \pi/\rho/a_1) = 54.77\% \quad (\text{PDG})$$

$$BF(\tau \rightarrow e/\mu) = 35.20\%$$

- One final lepton helps to remove huge QCD background.

➤ Collinear Approximation of τ^\pm decay.

K. Hagiwara, A. D. Martin and D. Zeppenfeld(1990)

➤ Reconstruction of scalar mass.

- $$\left(\frac{1}{x_l} - 1\right) P_T^l + \left(\frac{1}{x_h} - 1\right) P_T^h = P_T$$
 D. Rainwater, D. Zeppenfeld and K. Hagiwara(1998)
- $$P_{\tau^1} = \frac{P_l}{x_l} \qquad \qquad P_{\tau^2} = \frac{P_h}{x_h}$$
 Chung Kao, Duane A. Dicus, Rahul Malhotra and Yili Wang(2008)
- $$M_{\tau\tau} = \sqrt{(P_{\tau^1} + P_{\tau^2})^2}$$



- Total cross section of signal.

$$\sigma_{tot} = \int [f_b(\epsilon_1, \mu_F) \bar{f}_g(\epsilon_2, \mu_F) \rightarrow f_b(\epsilon_2, \mu_F) \bar{f}_g(\epsilon_1, \mu_F)] g \rightarrow b\phi \rightarrow b\tau^+\tau^-$$

$$\times 2 \times 2 BF(\tau \rightarrow l) BF(\tau \rightarrow j_\tau)$$

- Five flavor parton distribution function: CETQ6L1

- Factorization scale: $\mu_F = \frac{M_\phi}{4}$
- Renormalization scale: $\mu_R = \frac{M_\phi}{4}$

- Background processes.

- Drell-Yan processes:

$$q/\bar{q}g \rightarrow q/\bar{q}Z^*/\gamma^* \rightarrow q/\bar{q}\tau^+\tau^- \quad (q=u,d,s,c)$$

- To include the higher order correction, K factor is chosen to be 1.3.



- $t\bar{t}$ production:

$$q\bar{q}(gg) \rightarrow t\bar{t} \rightarrow b\bar{b} W^+ W^- \rightarrow b\bar{b} e^\pm(\mu^\pm) \tau^\mp + \cancel{E}$$

$$q\bar{q}(gg) \rightarrow t\bar{t} \rightarrow b\bar{b} W^+ W^- \rightarrow b\bar{b} \tau^+ \tau^- + \cancel{E}$$

$$q\bar{q}(gg) \rightarrow t\bar{t} \rightarrow b\bar{b} W^+ W^- \rightarrow b\bar{b} e^\pm(\mu^\pm) j_1 j_2 + \cancel{E}$$

$$q\bar{q}(gg) \rightarrow t\bar{t} \rightarrow b\bar{b} W^+ W^- \rightarrow b\bar{b} \tau^\pm j_1 j_2 + \cancel{E}$$

$$q\bar{q}(gg) \rightarrow t\bar{t} \rightarrow b\bar{b} W^+ W^- \rightarrow b\bar{b} \tau^\pm j_1 j_2 + \cancel{E}$$

- $j_1 j_2$ are quark and anti-quark pair from W decay.
- $K = 2.0$

- tW production:

$$b(\bar{b})g \rightarrow t(\bar{t})W^\mp \rightarrow b(\bar{b})W^\pm W^\mp \rightarrow b(\bar{b})e^\mp(\mu^\mp) \tau^\pm + \cancel{E}$$

$$b(\bar{b})g \rightarrow t(\bar{t})W^\mp \rightarrow b(\bar{b})W^\pm W^\mp \rightarrow b(\bar{b})\tau^\mp \tau^\pm + \cancel{E}$$

$$b(\bar{b})g \rightarrow t(\bar{t})W^\mp \rightarrow b(\bar{b})W^\pm W^\mp \rightarrow b(\bar{b})e^\mp(\mu^\mp) j_1 j_2 + \cancel{E}$$

$$b(\bar{b})g \rightarrow t(\bar{t})W^\mp \rightarrow b(\bar{b})W^\pm W^\mp \rightarrow b(\bar{b})\tau^\mp j_1 j_2 + \cancel{E}$$

- $K = 1.5$



- The production of $j_1 j_2 W$ and bjW are negligible.

➤ Acceptance cuts.

$$\sqrt{s} = 14 \text{ TeV}$$

$$\int L dt = 30 \text{ fb}^{-1}$$

$$P_T(b) > 15 \text{ GeV}$$

$$\cancel{E}_T > 20 \text{ GeV}$$

$$|M_{\tau\tau} - M_\phi| < 0.15 M_\phi \quad |M_{\tau\tau} - M_\phi| < 0.20 M_\phi$$

$$\sqrt{s} = 7 \text{ TeV}$$

$$\int L dt = 1 \text{ fb}^{-1}$$

$$P_T(b) > 15 \text{ GeV}$$

$$E_T > 20 \text{ GeV}$$

$$|M_{\tau\tau} - M_\phi| < 0.15 M_\phi$$

$$P_T(l) > 20 \text{ GeV} \quad P_T(j_\tau) > 40 \text{ GeV}$$

$$\eta(b) < 2.5 \quad \eta(l) < 2.5 \quad \eta(j_\tau) < 2.5$$

$$\Phi(l, j_\tau) < 170^\circ \quad \delta R(l, j_\tau) > 0.3 \quad M(l, \cancel{E}_T) < 30 \text{ GeV}$$

- One more set of cuts is required by the physical meaning of energy fraction.

$$0 < x_l < 1 \quad 0 < x_h < 1$$



- Tagging efficiency and mistagging efficiency.

$$\sqrt{s} = 14 \text{TeV}$$

$$\int L dt = 30 \text{fb}^{-1}$$

$$\varepsilon_b = 60\%$$

$$\sqrt{s} = 7 \text{TeV}$$

$$\int L dt = 300 \text{fb}^{-1}$$

$$\varepsilon_b = 50\%$$

$$\int L dt = 1 \text{fb}^{-1}$$

$$\varepsilon_b = 50\%$$

$$\varepsilon_{l_\tau} = 26\%$$

$$P_{g,u,d,s \rightarrow b} = 1\% \quad P_{c \rightarrow b} = 10\% \quad P_{u,d,c,s \rightarrow j_\tau} = 1/400 \quad P_{b \rightarrow j_\tau} = 1/600$$

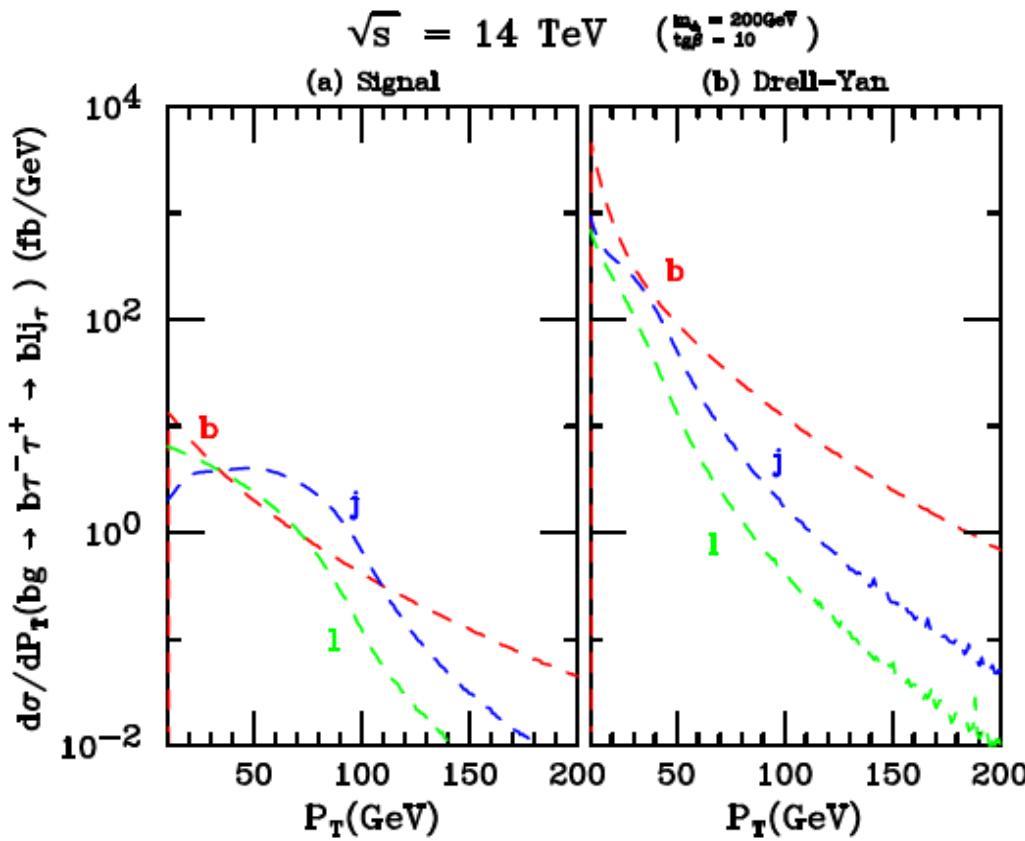
- Criterion for the observability of some signal .

$$\sigma_s > \frac{N^2}{L} \left(+ 2\sqrt{L\sigma_b} / N \right)$$

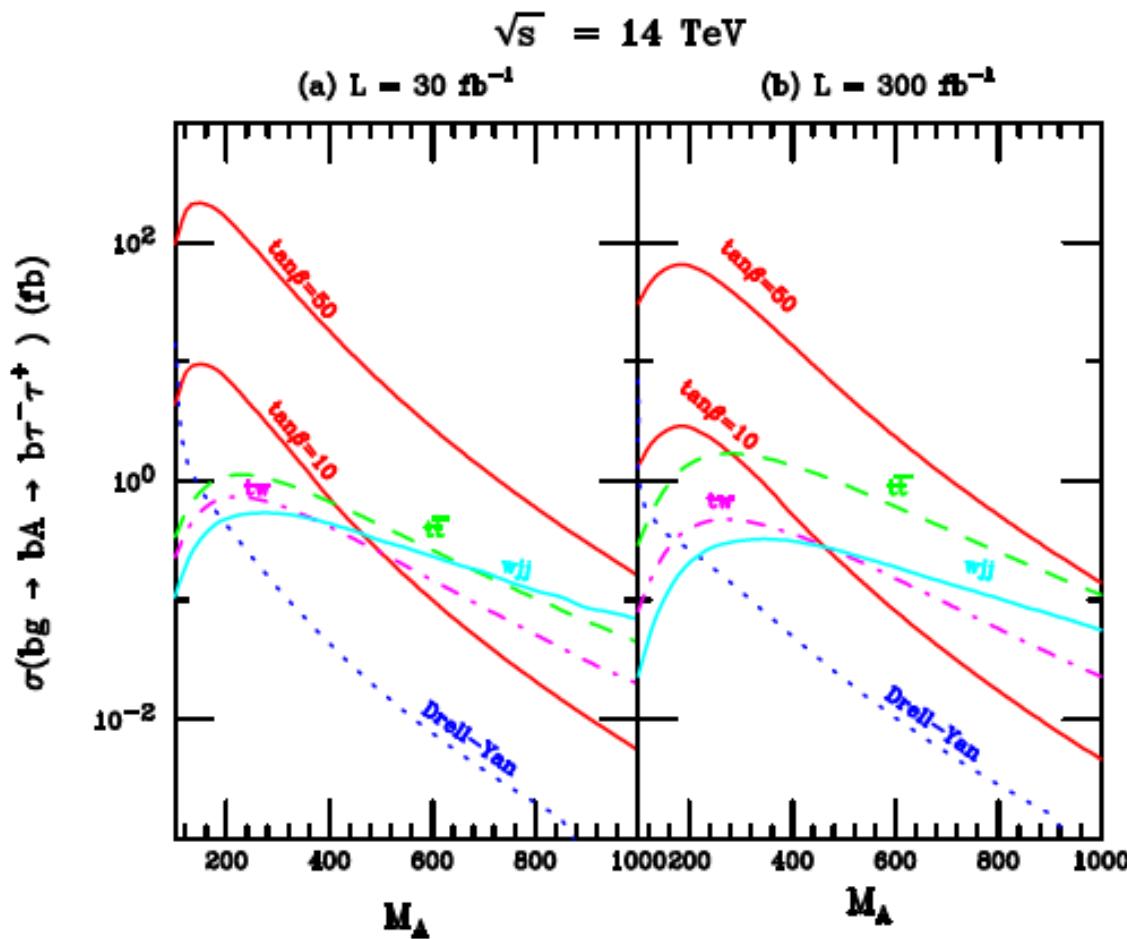
- $N = 2.5$ corresponds to 5σ .



- P_T distribution without any cuts.



➤ $\sigma \sim m_A$



➤ Signal significance.

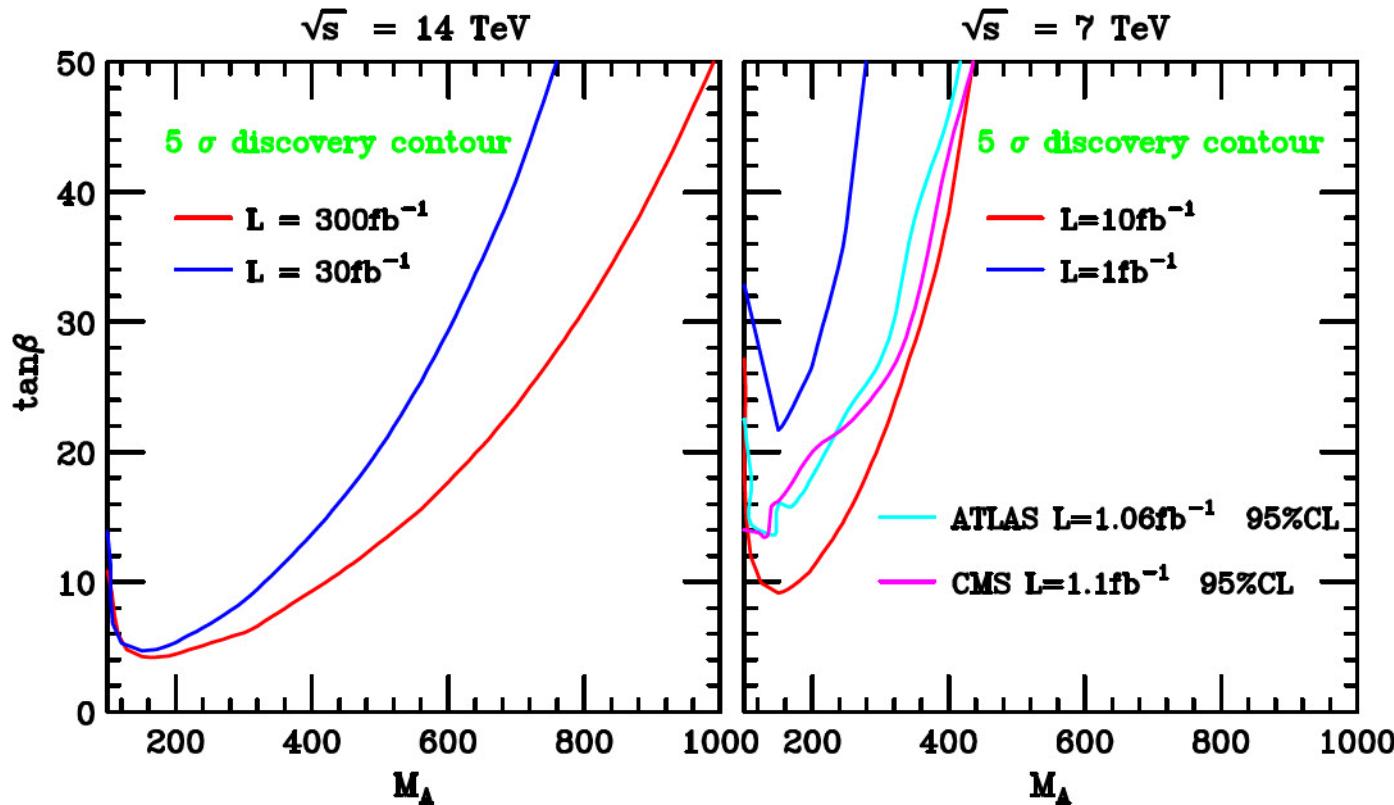
TABLE I: MSSM Higgs Production at $\sqrt{s} = 14\text{TeV}$ and $\mathcal{L} = 30\text{fb}^{-1}$

$M_A(\text{GeV})$	100	200	400	800
$\sigma_s(\tan\beta = 10)$	4.38	7.28	7.14×10^{-1}	2.09×10^{-2}
$\sigma_s(\tan\beta = 50)$	9.75×10^1	1.61×10^2	1.81×10^1	6.02×10^{-1}
$\sigma_s(\text{Drell-Yan})$	1.41×10^1	4.27×10^{-1}	4.32×10^{-2}	1.99×10^{-3}
$\sigma_s(b\bar{b}W^+W^-)$	3.41×10^{-1}	1.10	6.75×10^{-1}	1.04×10^{-1}
$\sigma_s(bW^+W^-)$	2.30×10^{-1}	7.25×10^{-1}	4.13×10^{-1}	5.16×10^{-2}
$\sigma_s(Wjj)$	1.07×10^{-1}	4.83×10^{-1}	4.40×10^{-1}	1.21×10^{-1}
$N_{ss}(\tan\beta = 10)$	6.24	24.1	3.12	0.217
$N_{ss}(\tan\beta = 50)$	139	533	79.1	6.25

TABLE II: MSSM Higgs Production at $\sqrt{s} = 14\text{TeV}$ and $\mathcal{L} = 300\text{fb}^{-1}$

$M_A(\text{GeV})$	100	200	400	800
$\sigma_s(\tan\beta = 10)$	1.34	2.83	5.21×10^{-1}	1.74×10^{-2}
$\sigma_s(\tan\beta = 50)$	3.08×10^1	6.40×10^1	1.37×10^1	5.19×10^{-1}
$\sigma_s(\text{Drell-Yan})$	6.97	2.64×10^{-1}	4.98×10^{-2}	2.84×10^{-3}
$\sigma_s(b\bar{b}W^+W^-)$	2.91×10^{-1}	1.39	1.35	2.55×10^{-1}
$\sigma_s(bW^+W^-)$	8.09×10^{-2}	4.00×10^{-1}	3.61×10^{-1}	5.78×10^{-2}
$\sigma_s(Wjj)$	2.31×10^{-2}	2.06×10^{-1}	3.13×10^{-1}	1.04×10^{-1}
$N_{ss}(\tan\beta = 10)$	8.55	32.6	6.27	0.465
$N_{ss}(\tan\beta = 50)$	197	737	165	13.9

- 5σ discovery contour.



Several Experimental Constraints

➤ $B_s^0 \rightarrow \mu^+ \mu^-$

$$BF(B_s^0 \rightarrow \mu^+ \mu^-)^{\text{SM}} = (3.2 \pm 0.2) \times 10^{-9}$$

A. J. Buras arXiv:1012.1447v2

$$BF(B_s^0 \rightarrow \mu^+ \mu^-)^{\text{EXP}} < 1.08 \times 10^{-8} \quad 95\% \text{ C.L.}$$

LHCb-CONF-2011-047

➤ $b \rightarrow s \gamma$

$$\begin{aligned} BF(b \rightarrow s \gamma)^{\text{SM}} &= (2.98 \pm 0.26) \times 10^{-4} \\ &= (3.15 \pm 0.23) \times 10^{-4} \end{aligned}$$

Becher & Neubert (2007)

Misiak et al. (2007)

$$BF(b \rightarrow X_s \gamma)^{\text{EXP}} = (3.45 \pm 0.15 \pm 0.40) \times 10^{-4}$$

Wenfeng Wang (2011)

➤ Δa_μ

$$a_\mu^{\text{EXP}} = 116592089(6.3) \times 10^{-10}$$

PDG (2010)

$$a_\mu^{\text{SM}} = 116591830(5.1) \times 10^{-10}$$

T. Teubner (2010)

$$\Delta a_\mu \equiv a_\mu^{\text{EXP}} - a_\mu^{\text{SM}} = (25.9 \pm 8.1) \times 10^{-10}$$

Gi-Chol Cho, Kaoru Hagiwara, Yu Matsumoto
and Daisuke Nomura (2011)

mSUGRA Higgs Discovery Potential

- SUSY breaking happens in a hidden sector, and is mediated to visible sector by a messenger, gravity.
- mSUGRA/CMSSM assumes unified scalar mass m_0 , fermionic mass $m_{1/2}$ and trilinear coupling A_0 at GUT scale. These are free inputs.
- In addition, two more parameters are defined at low-energy scale, $\tan \beta$ and $sign(\mu)$.
- RGEs evolve from GUT scale down to EW scale, and then generate particle spectrum at EW scale.

Isajet 7.81 (H. Baer, F.E. Paige, S.D. Protopopescu, X. Tata)

- The following parameter space will be scanned.

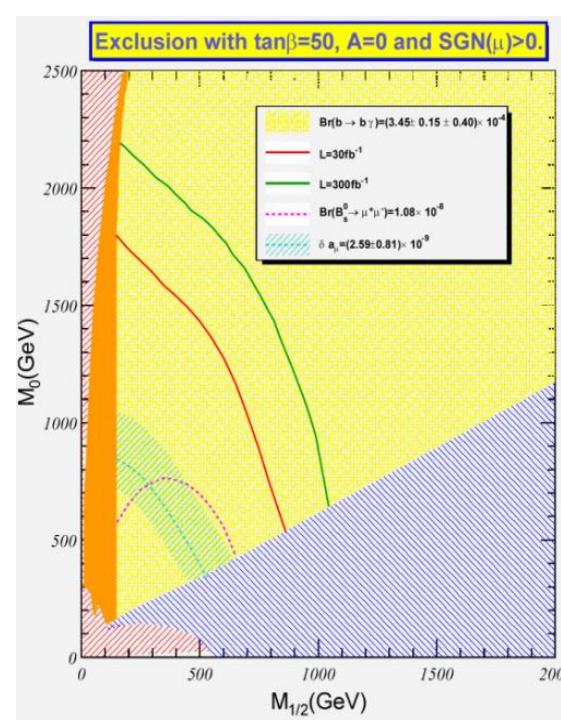
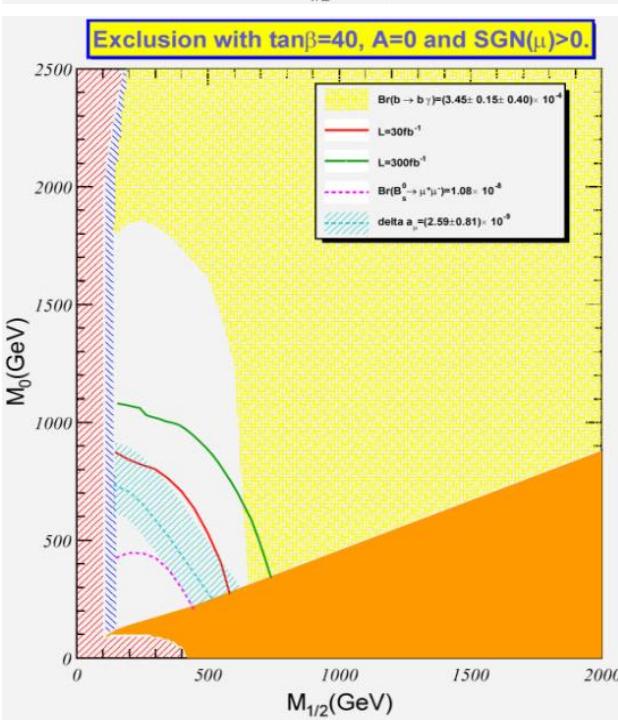
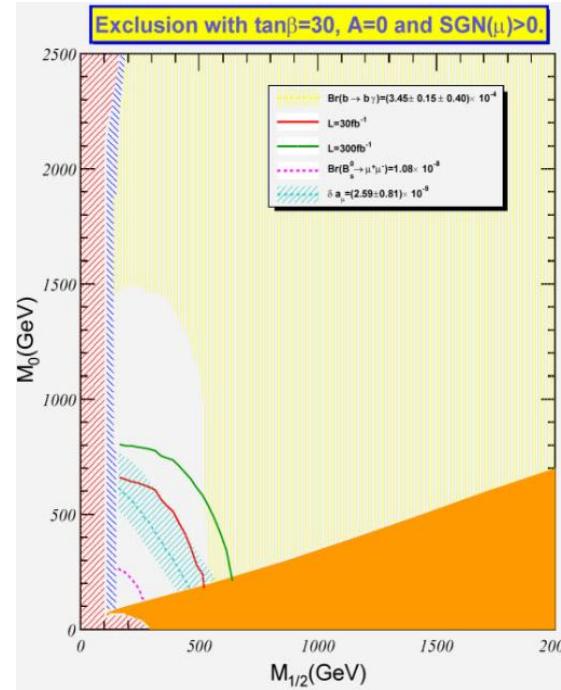
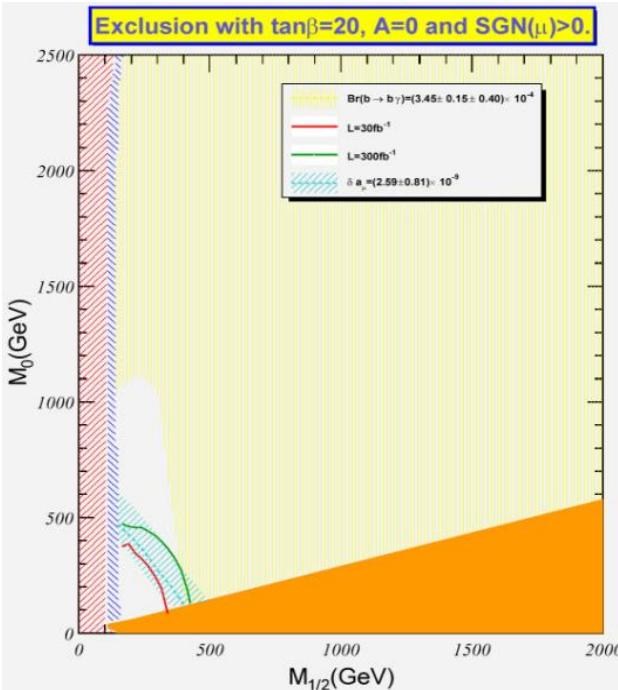
$$m_{1/2} \in [0, 2000] \text{ GeV}, \quad m_0 \in [0, 25000] \text{ GeV}, \quad \tan \beta \in [20, 50]$$

$$A_0 = 0 \text{ GeV}, \quad sign(\mu) = 1$$

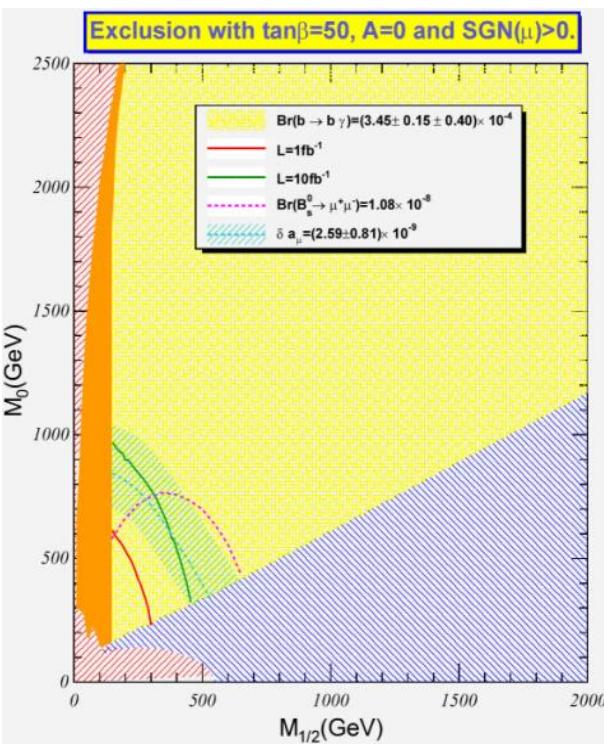
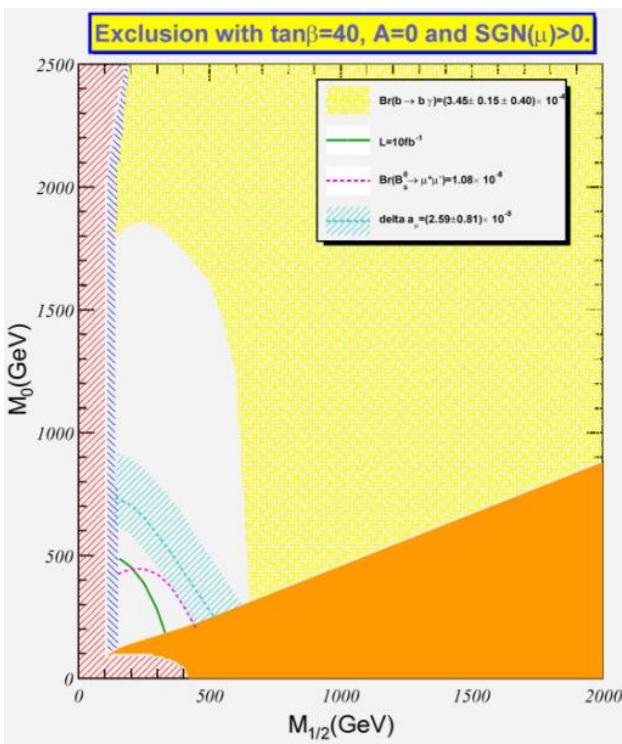
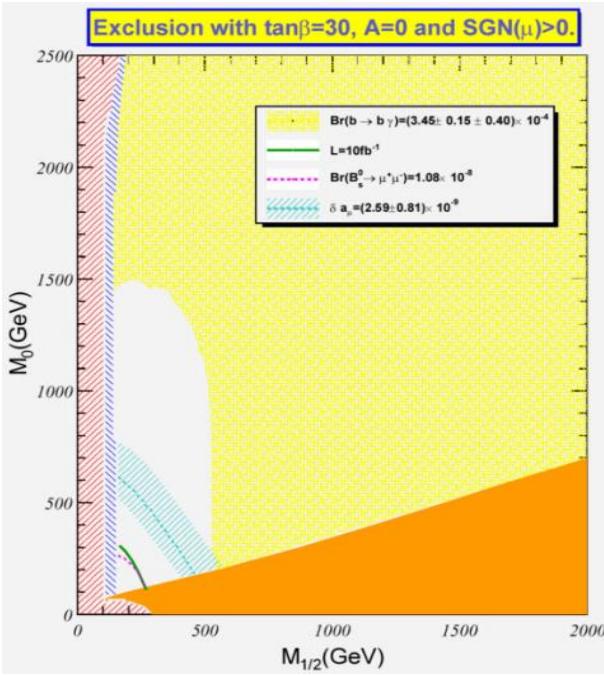
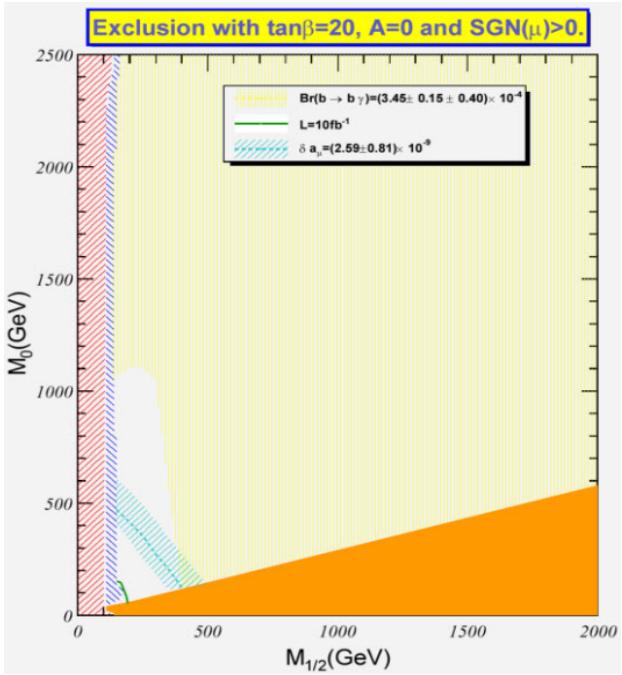
$$m_{top} = 173.1 \text{ GeV}$$



14TeV



7TeV



mAMSB Higgs Discovery Potential

- SUSY breaking happens in a separate brane, and is mediated to visible sector by super-Weyl anomaly.
- The parameter space is formed by four parameters.

$$\{m_{3/2}, m_0, \tan\beta, sign(\mu)\}$$

- The scanned parameter space

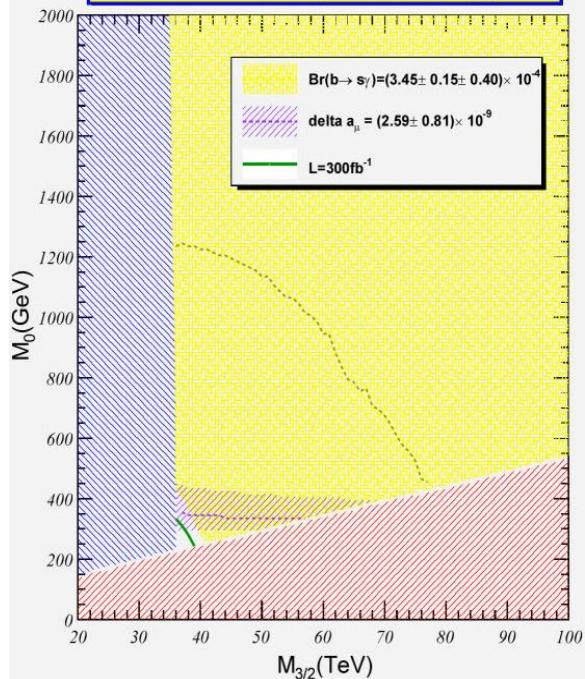
$$m_{3/2} \in [20000, 100000], \quad m \in [0, 2000], \quad \tan\beta \in [20, 50]$$

$$sign(\mu) = 1$$

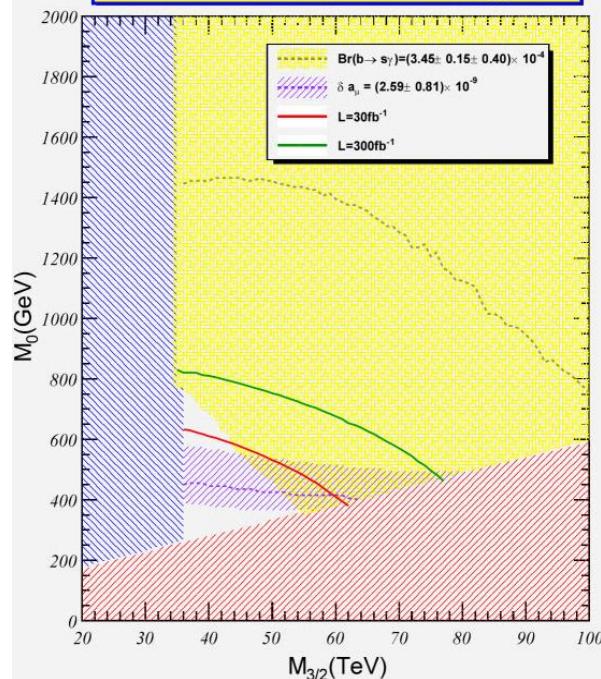
$$m_{top} = 173.1 GeV$$



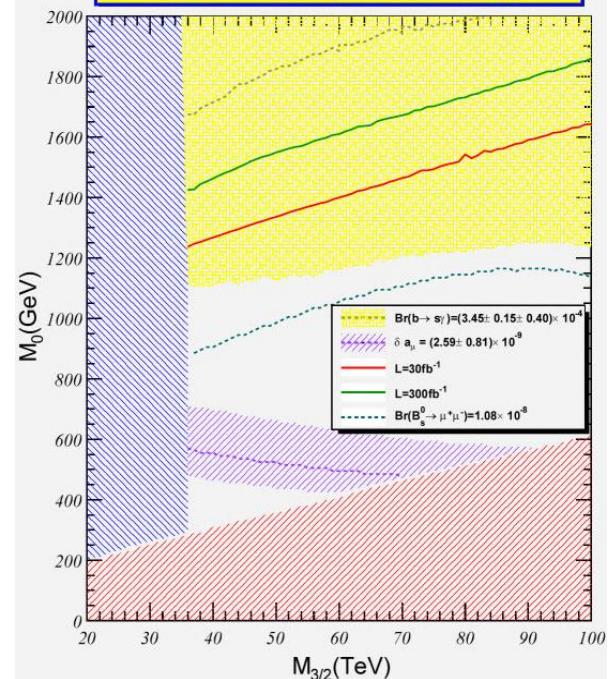
Exclusion with $\tan\beta=20$ and $\text{SGN}(\mu)>0.$



Exclusion with $\tan\beta=30$ and $\text{SGN}(\mu)>0.$



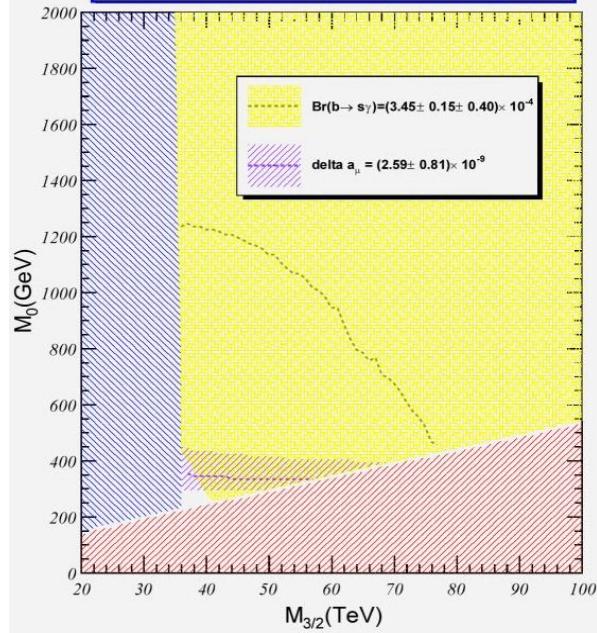
Exclusion with $\tan\beta=40$ and $\text{SGN}(\mu)>0.$



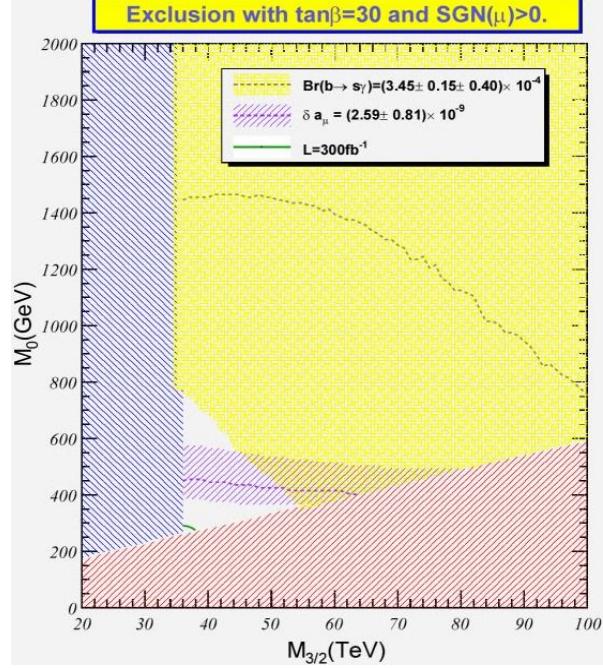
$14TeV$



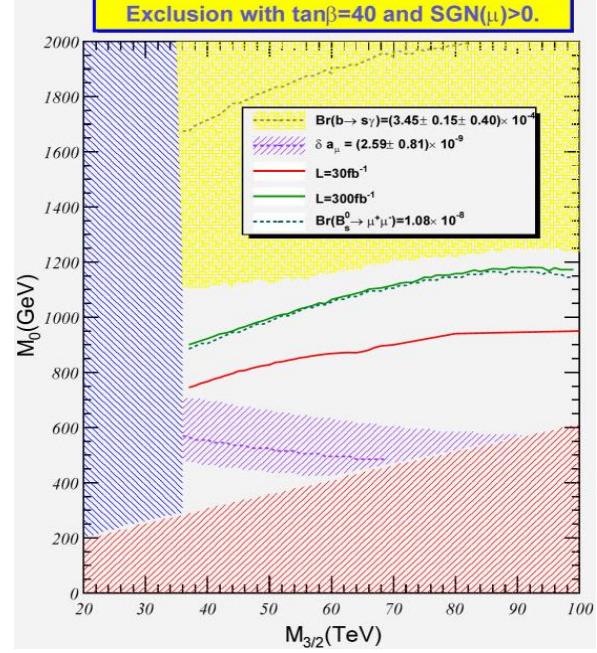
Exclusion with $\tan\beta=20$ and $SGN(\mu)>0$.



Exclusion with $\tan\beta=30$ and $SGN(\mu)>0$.



Exclusion with $\tan\beta=40$ and $SGN(\mu)>0$.



7 TeV



Conclusion

- It's promising to discover Higgs by tau lepton pairs, even with Higgs mass up to 1 TeV at LHC.
- In mSUGRA, the model with high $\tan \beta$ favors the discovery of Higgs.
- In mAMSB, the model with intermediate $\tan \beta$ favors the discovery of Higgs.
- GMSB and related channels will be added.

